

VI Amazonian Symposium on Physics

18th-22nd November 2024
Federal University of Pará

OC12:

Gravitational bremsstrahlung and the Fulling-Davies-Unruh
thermal bath

Belém -Pará - Brazil

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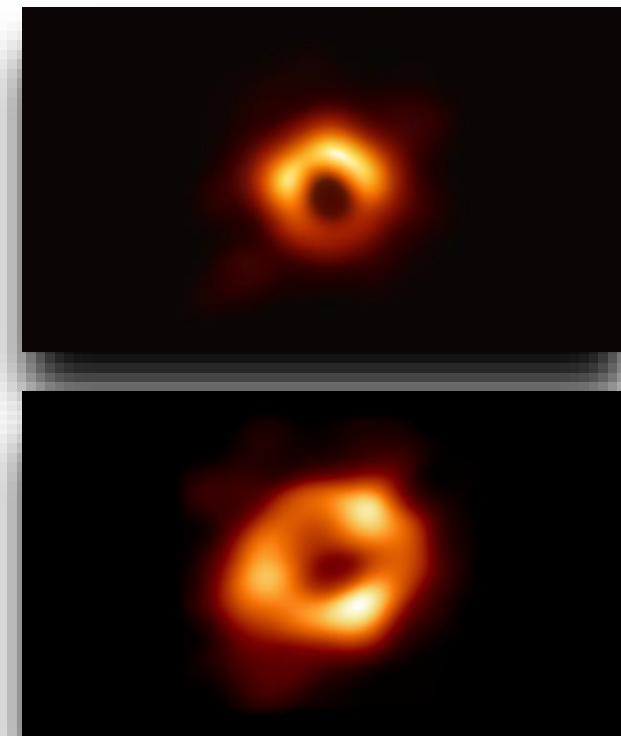
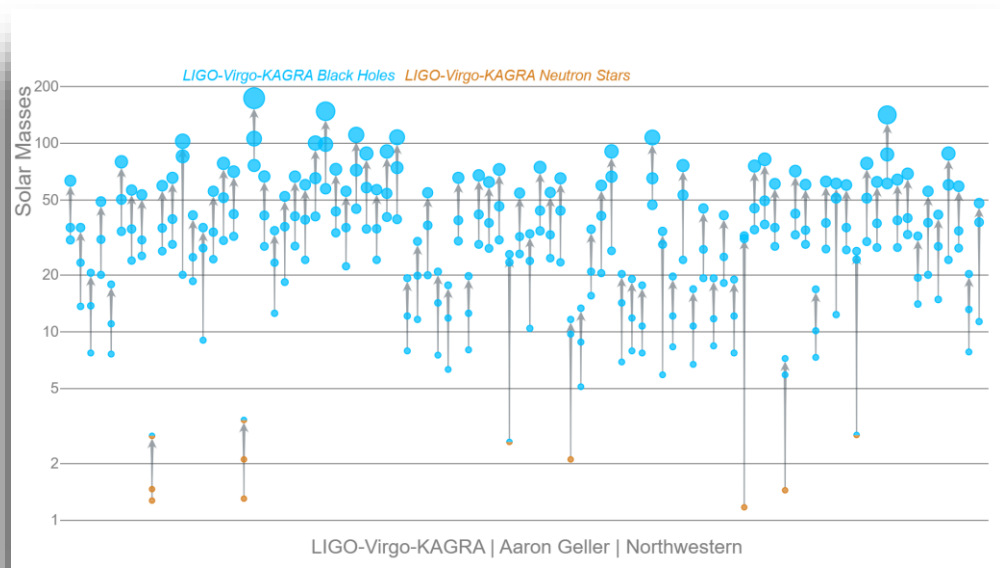
Belém-Pará

Thursday the 21st of November, 2024



■ LIGO-Virgo-KAGRA

■ EHT



GWs catalogue and the first shadows of BHs. Top: M87* BH. Bottom: Sgr A* BH.

Strong evidence for the existence of BHs: LIGO-Virgo-KAGRA + EHT

Graphic representation of the masses of announced GW detections in a Stellar Graveyard

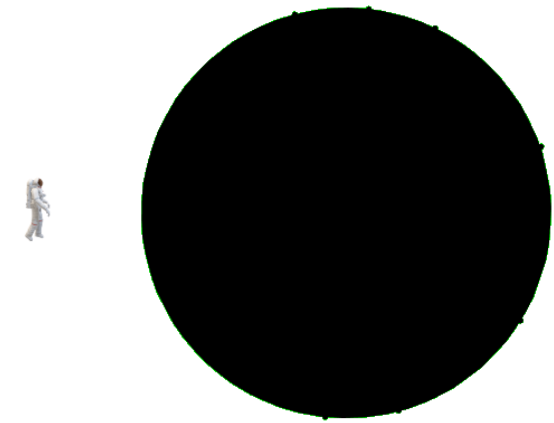


- Quantum gravity?



- Prime results: GR 🤝 QFT
 - Hawking effect

$$T_H = \frac{\hbar c^3}{8\pi k_B G M}$$



Pictorial representation of Hawking radiation.



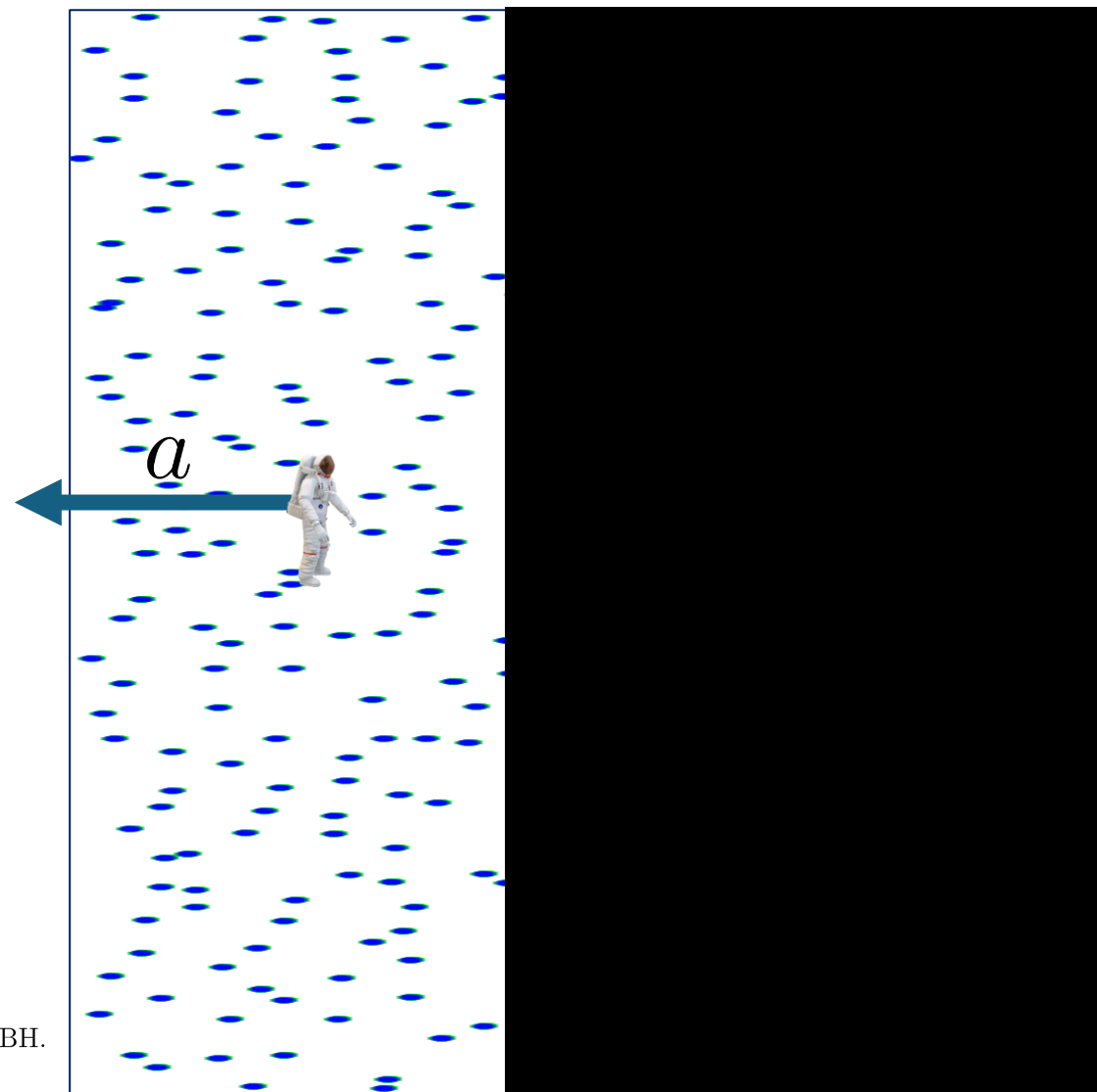
- Prime results: GR 🤝 QFT

- Hawking effect:

$$T_H = \frac{\hbar\kappa}{2\pi ck_B}$$

- Unruh effect:

$$T_U = \frac{\hbar a}{2\pi ck_B}$$



Representation of Hawking radiation in the vicinity of the BH.



The Unruh effect and its applications

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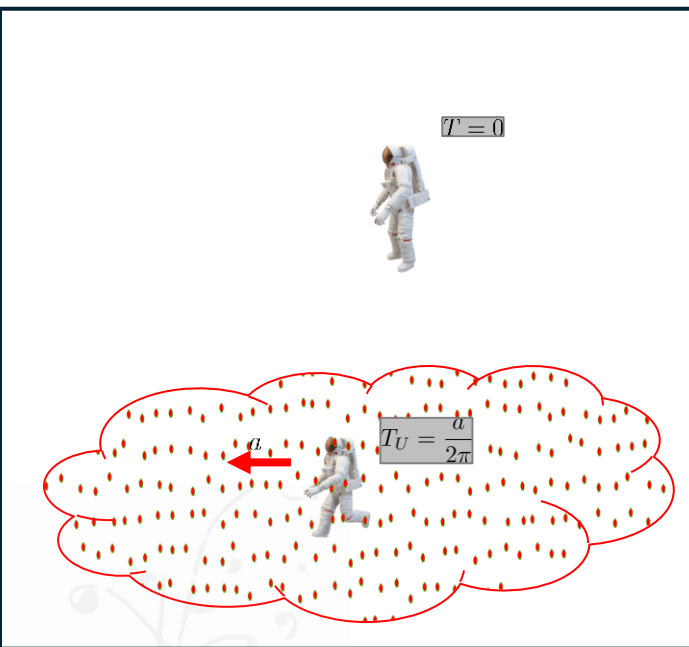
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$$\langle 0_M | \hat{N}_M | 0_M \rangle = 0$$



A pictorial representation of the Unruh effect.

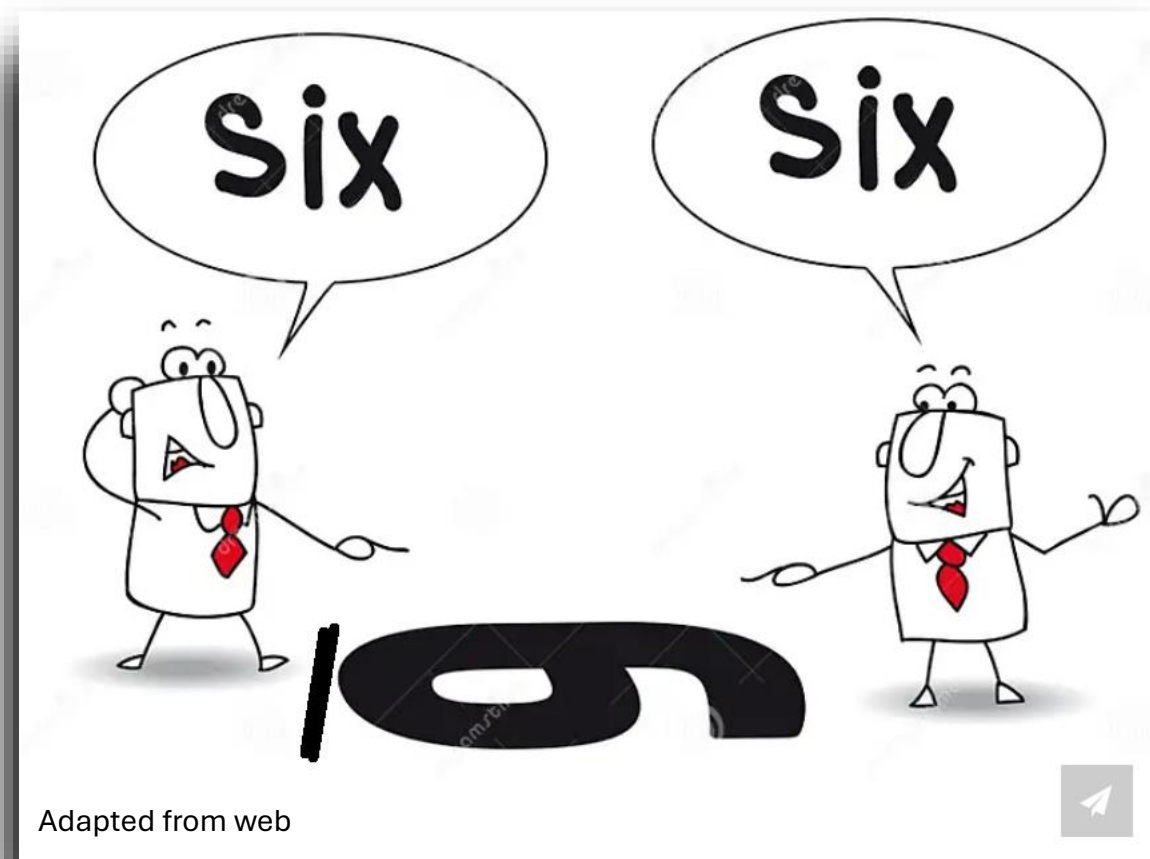
Nº of Rindler particles in $|0_M\rangle$:

$$n(\omega) = \frac{1}{e^{\frac{2\pi\omega}{a}} - 1}$$

This term first appeared in Bose, S. N. Z. Physik 26, 178–181 (1924);

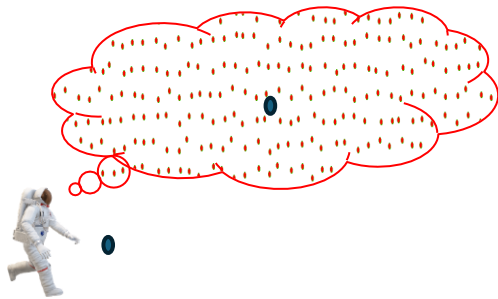
$$\langle 0_M | \hat{N}_R | 0_M \rangle = n(\omega)$$



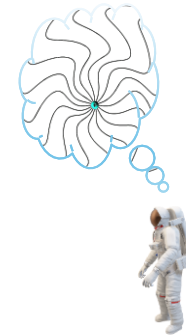


□ Radiation in flat spacetime

Accelerated observer:



Inertial observer:

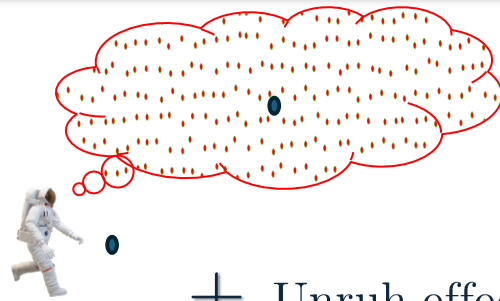


The response to the quantum field must be equal!

+ Unruh effect:

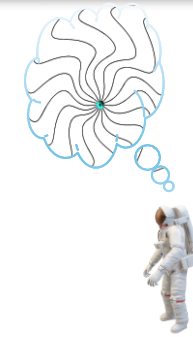
$$T_U = \frac{\hbar a}{2\pi c k_B}$$





+ Unruh effect:

$$T_U = \frac{\hbar a}{2\pi c k_B}$$



- All observers must agree with the response rate:

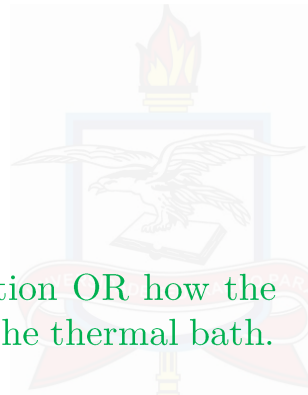
$$R_{\text{em}}^{\lambda \mathbf{k}_\perp} = \int_0^\infty d\omega \frac{|A_{\text{em}}^{\lambda \mathbf{k}_\perp}|^2}{T}$$

$$\mathcal{A}_{\text{em}}^{(\omega, \mathbf{k}_\perp)} = \langle 0_{\text{F}} | \hat{a}_{(\omega, \mathbf{k}_\perp)} \hat{S}_{\text{int}} | 0_{\text{F}} \rangle$$

- Interaction term describing the one-graviton processes:

$$\hat{S}_{\text{int}} = \sqrt{8\pi G} \int T^{\mu\nu} \hat{h}_{\mu\nu} \sqrt{-g} d^4x$$

We are analyzing gravitational radiation OR how the source interacts with gravitons from the thermal bath.



Theoretical description of bremsstrahlung:

Spontaneous emission

$$R_{em}^{\lambda \mathbf{k}_{\perp}} = \int_0^{\infty} d\omega \frac{|A_{em}^{\lambda \mathbf{k}_{\perp}}|^2}{T} = 0$$

Absorption/induced emission

$$R^{\lambda \mathbf{k}} = \int_0^{\infty} d\omega \frac{|A_{em}^{\lambda \mathbf{k}_{\perp}}|^2}{T} n(\omega) \neq 0$$

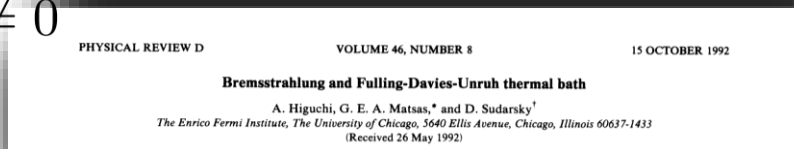
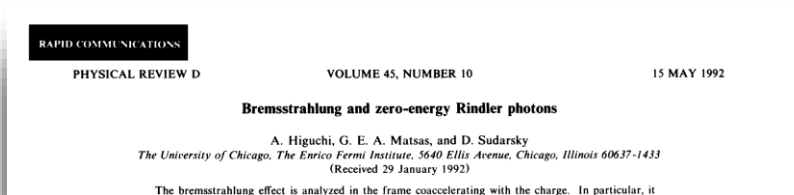
Nº of Rindler photons with zero energy

$$n(\omega) = \frac{1}{e^{\frac{2\pi\omega}{a}} - 1}$$

Polarization

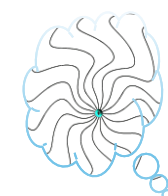
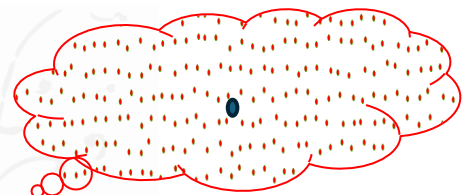
Transverse momentum

$$R_{em}^{\lambda \mathbf{k}_{\perp}} = \int_0^{\infty} d\omega \frac{|A_{em}^{\lambda \mathbf{k}_{\perp}}|^2}{T}$$



$$R_{k_{\perp}}^{\text{tot}} = \frac{q^2}{4\pi^3 a} \left| K_1 \left(\frac{k_{\perp}}{a} \right) \right|^2$$

A pictorial representation of the bremsstrahlung.



- Minkowski coordinates:

$$(t, z, x, y)$$

- *Right* Rindler wedge:

$$t = a^{-1} e^{a\xi} \sinh a\eta,$$

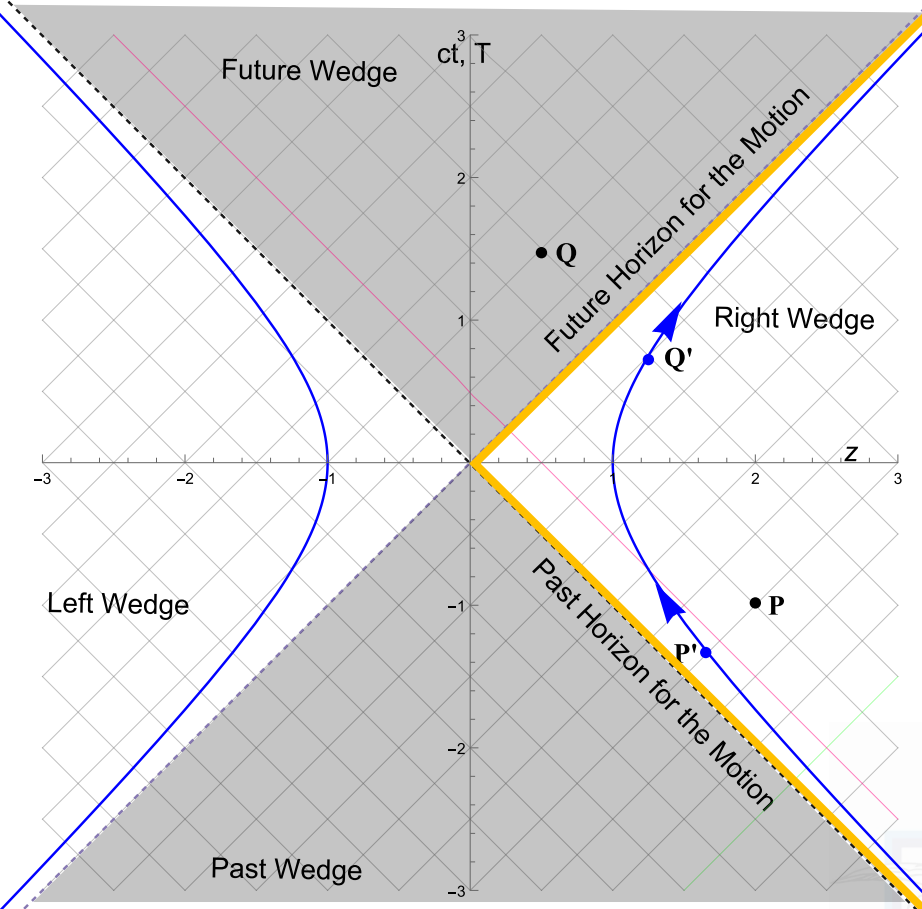
$$z = a^{-1} e^{a\xi} \cosh a\eta,$$

where $z > |t|$ and $-\infty < \eta, \xi < \infty$.

- Accelerated-observer spacetime:

$$d\tau^2 = e^{2a\xi} (d\eta^2 - d\xi^2) - dx^2 - dy^2$$

Proper acceleration $\alpha := ae^{-a\xi}$.
 Proper time $\tau = e^{a\xi}\eta$.



Minkowski spacetime diagram depicting hyperbolic trajectories and the four Rindler wedges.

- Inertial observer (Minkowski coordinates):

$$T^{tt} = \mu [\cosh^2 a\eta \delta(\xi) + a\theta(\xi)] \delta^{(2)}(\mathbf{x}_\perp),$$

$$T^{tz} = \mu \sinh a\eta \cosh a\eta \delta(\xi) \delta^{(2)}(\mathbf{x}_\perp),$$

$$T^{zz} = \mu [\sinh^2 a\eta \delta(\xi) - a\theta(\xi)] \delta^{(2)}(\mathbf{x}_\perp).$$

- Conservation equation:

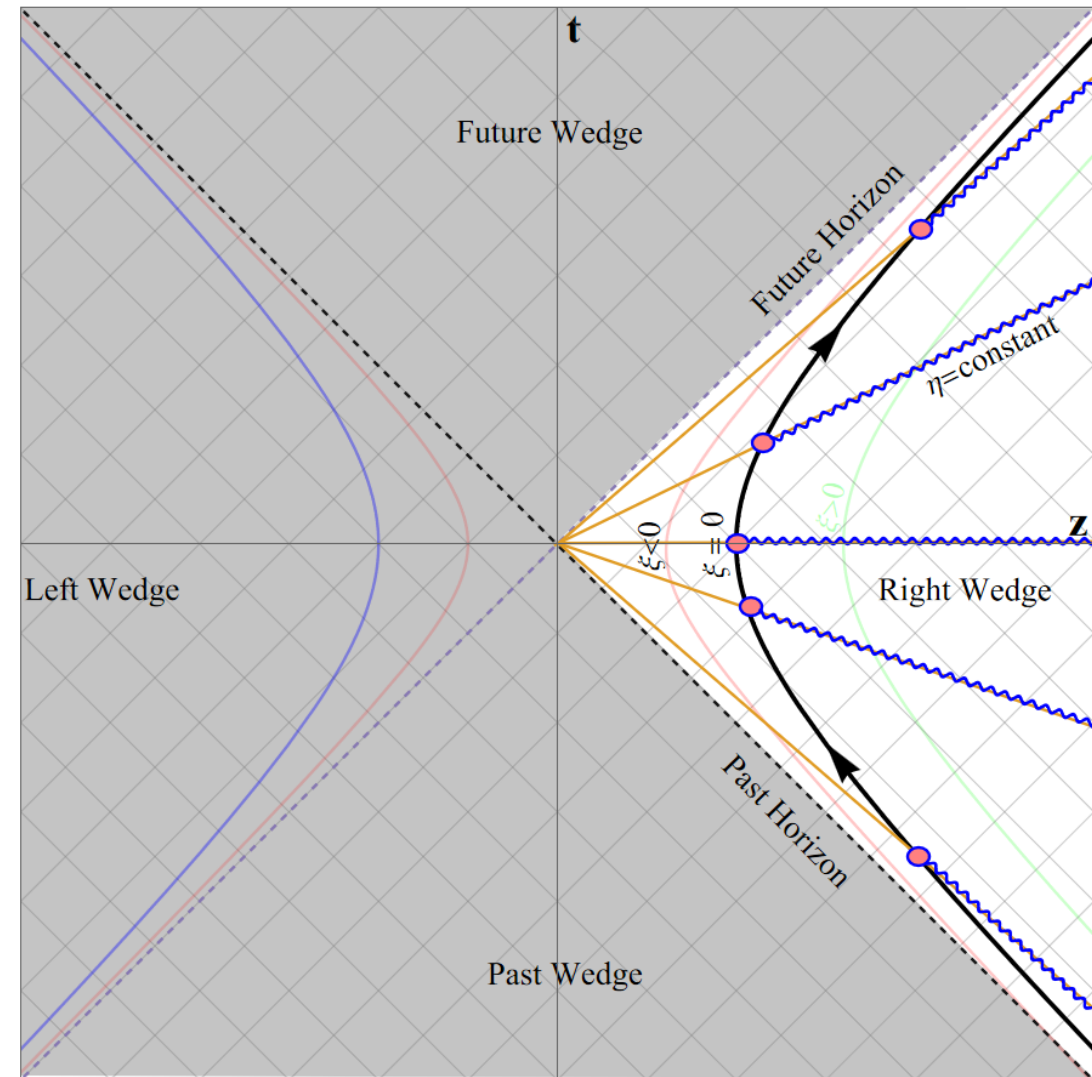
$$\nabla_\mu T^{\mu\nu} = 0$$

- Accelerated observer (Rindler coordinates):

$$T^{\eta\eta} = \mu \left[\delta(\xi) + (1 - \beta) a e^{-(2+\beta)a\xi} \theta(\xi) \right] \delta^{(2)}(\mathbf{x}_\perp),$$

$$T^{\xi\xi} = -\mu a e^{-(2+\beta)a\xi} \theta(\xi) \delta^{(2)}(\mathbf{x}_\perp).$$

Minkowski spacetime with the stress-energy depicted in the right Rindler wedge.



- The total emission probability:

$$\mathcal{P}^M = \int \frac{d^3\mathbf{k}}{(2\pi)^3 2k} \overline{\mathcal{T}^{\mu\nu}(\mathbf{k})} \sum_{\lambda} \epsilon_{\mu\nu}^{\lambda} \epsilon_{\alpha\beta}^{\lambda} \mathcal{T}^{\alpha\beta}(\mathbf{k})$$

Where $\epsilon_{\mu\nu}^{\lambda}$ is the Minkowski polarization matrix:

Considering $k^{\mu} = (k, 0, 0, k)$

$$\epsilon_{11}^{\lambda_1} = -\epsilon_{22}^{\lambda_1} = \epsilon_{12}^{\lambda_2} = \epsilon_{21}^{\lambda_2} = \frac{1}{\sqrt{2}}$$

- With fixed transverse momentum:

$$\begin{aligned} \mathcal{P}_{\mathbf{k}_{\perp}}^M &= \int \frac{dk_z}{(2\pi)^3 2k} \left(\overline{\mathcal{T}^{\mu\nu}} \mathcal{T}_{\mu\nu} - \frac{1}{2} \overline{\mathcal{T}} \mathcal{T} \right) \\ &= \int \frac{dk_z}{(2\pi)^3 2k} \left(\frac{1}{2} |\mathcal{T}^{tt} + \mathcal{T}^{zz}|^2 - 2 |\mathcal{T}^{tz}|^2 \right), \end{aligned}$$

$$\mathcal{T}^{\mu\nu}(\mathbf{k}) = \int e^{i(kt - \mathbf{k} \cdot \mathbf{x})} T^{\mu\nu}(x) d^4x$$



- The differential response rate:

$$\begin{aligned} \mathcal{R}_{\mathbf{k}_\perp}^M &= \frac{\mathcal{P}_{\mathbf{k}_\perp}^M}{T_0} \\ &= \mu^2 \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} \cosh 2a\sigma \exp \left\{ \frac{2i}{a} \left[k' \sinh \left(\frac{a}{2} \sigma \right) \right] \right\} d\sigma \right] \frac{dk'_z}{4k'(2\pi)^3} \end{aligned}$$

Convergence factor by letting $\sigma \mapsto \sigma + 2i\epsilon$

$$\mathcal{R}_{\mathbf{k}_\perp}^M = \mu^2 \int_{-\infty}^{\infty} \left\{ \int_{-\infty}^{\infty} \cosh 2a\sigma \exp \left[\frac{ik'}{a} \left(e^{ia\epsilon} e^{a\sigma/2} - e^{-ia\epsilon} e^{-a\sigma/2} \right) \right] d\sigma \right\} \frac{dk'_z}{4k'(2\pi)^3}$$

$$s_\pm = \frac{k' + k'_z}{k_\perp} e^{\pm a\sigma/2} \int_0^\infty ds s^{\nu-1} \exp \left[\frac{i\mu}{2} \left(s - \frac{\beta^2}{s} \right) \right] = 2\beta^\nu e^{i\nu\pi/2} K_\nu(\beta\mu)$$

- The differential response rate:

$$\mathcal{R}_{\mathbf{k}_\perp}^M = \frac{\mu^2}{8\pi^3 a} \left| K_2 \left(\frac{k_\perp}{a} \right) \right|^2$$



- GWs in Rindler spacetime:

$$h_{\mu\nu}^{(\omega, \mathbf{k}_\perp)}(\eta, \xi, \mathbf{x}_\perp) = H_{\mu\nu}^{(\omega, k_\perp)}(\xi) e^{-i\omega\eta + i\mathbf{k}_\perp \cdot \mathbf{x}_\perp}$$

Positive frequency (∂_η)

Parity: Even Odd.

Even-parity modes:

$$\tilde{h}_{\mu\nu}^{R,e} = \begin{pmatrix} -\frac{\kappa^2(\xi)}{\kappa^2} \varphi^{R,e} - \frac{1}{\kappa^2} (\omega^2 \varphi^{R,e} + 2a \partial_\xi \varphi^{R,e}) & -\frac{2i\omega}{\kappa^2} (\partial_\xi \varphi^{R,e} - a \varphi^{R,e}) & 0 & 0 \\ * & -e^{2a\xi} \varphi^{R,e} - \frac{2}{\kappa^2} (a \partial_\xi \varphi^{R,e} - \partial_\xi^2 \varphi^{R,e}) & 0 & 0 \\ * & * & \varphi^{R,e} & 0 \\ * & * & * & \varphi^{R,e} \end{pmatrix}$$

$$\tilde{h}_{\mu\nu}^{R,o} = \begin{pmatrix} 0 & 0 & -\frac{ik_y}{\kappa^2} \partial_\xi \varphi^{R,o} & \frac{ik_x}{\kappa^2} \partial_\xi \varphi^{R,o} \\ * & 0 & -\frac{\omega k_x}{\kappa^2} \varphi^{R,o} & \frac{\omega k_y}{\kappa^2} \varphi^{R,o} \\ * & * & 0 & 0 \\ * & * & * & 0 \end{pmatrix}$$

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Gravitational waves in Kasner spacetimes and Rindler wedges in Regge-Wheeler gauge: Formulation of Unruh effect

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$$\nabla_\alpha \nabla^\alpha \varphi^{(\omega, \mathbf{k}_\perp)} = 0$$

$$\varphi^{(\omega, \mathbf{k}_\perp)}(\eta, \xi, \mathbf{x}_\perp) = \phi^{(\omega, k_\perp)}(\xi) e^{-i\omega\eta + i\mathbf{k}_\perp \cdot \mathbf{x}_\perp}$$

Normalized scalar modes:

$$\phi^{(\omega, k_\perp)}(\xi) = \sqrt{\frac{\sinh(\pi\omega/a)}{4\pi^4 a}} K_{i\omega/a} \left(\frac{k_\perp e^{a\xi}}{a} \right)$$

- The differential response rate:

$$\mathcal{R}_{\mathbf{k}_\perp}^R = \frac{\mu^2}{8\pi^3 a} \left| K_2 \left(\frac{k_\perp}{a} \right) \right|^2$$

- Explicitly, we have:

$$\begin{aligned} H_{\eta\eta}^{(\omega, k_\perp)}(\xi) &= H_{\xi\xi}^{(\omega, k_\perp)}(\xi) \\ &= \frac{1}{\sqrt{2}} \left[-e^{2a\xi} + \frac{2}{k_\perp^2} (\omega^2 + a\partial_\xi) \right] \phi^{(\omega, k_\perp)}(\xi), \end{aligned}$$

$$H_{\eta\xi}^{(\omega, k_\perp)}(\xi) = \frac{\sqrt{2}i\omega}{k_\perp^2} (\partial_\xi - a) \phi^{(\omega, k_\perp)}(\xi),$$

$$H_{ij}^{(\omega, k_\perp)}(\xi) = -\frac{\delta_{ij}}{\sqrt{2}} \phi^{(\omega, k_\perp)}(\xi),$$



- The consistent description of physics (in $|0_M\rangle$) in the accelerated frame requires the Unruh effect.
- In particular, we have fundamentally described the bremsstrahlung in relation to the Unruh effect.



□ Acknowledgments



- Prof. Crispino
- Prof. Higuchi



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