



# *Compact binaries in astrophysical environments*

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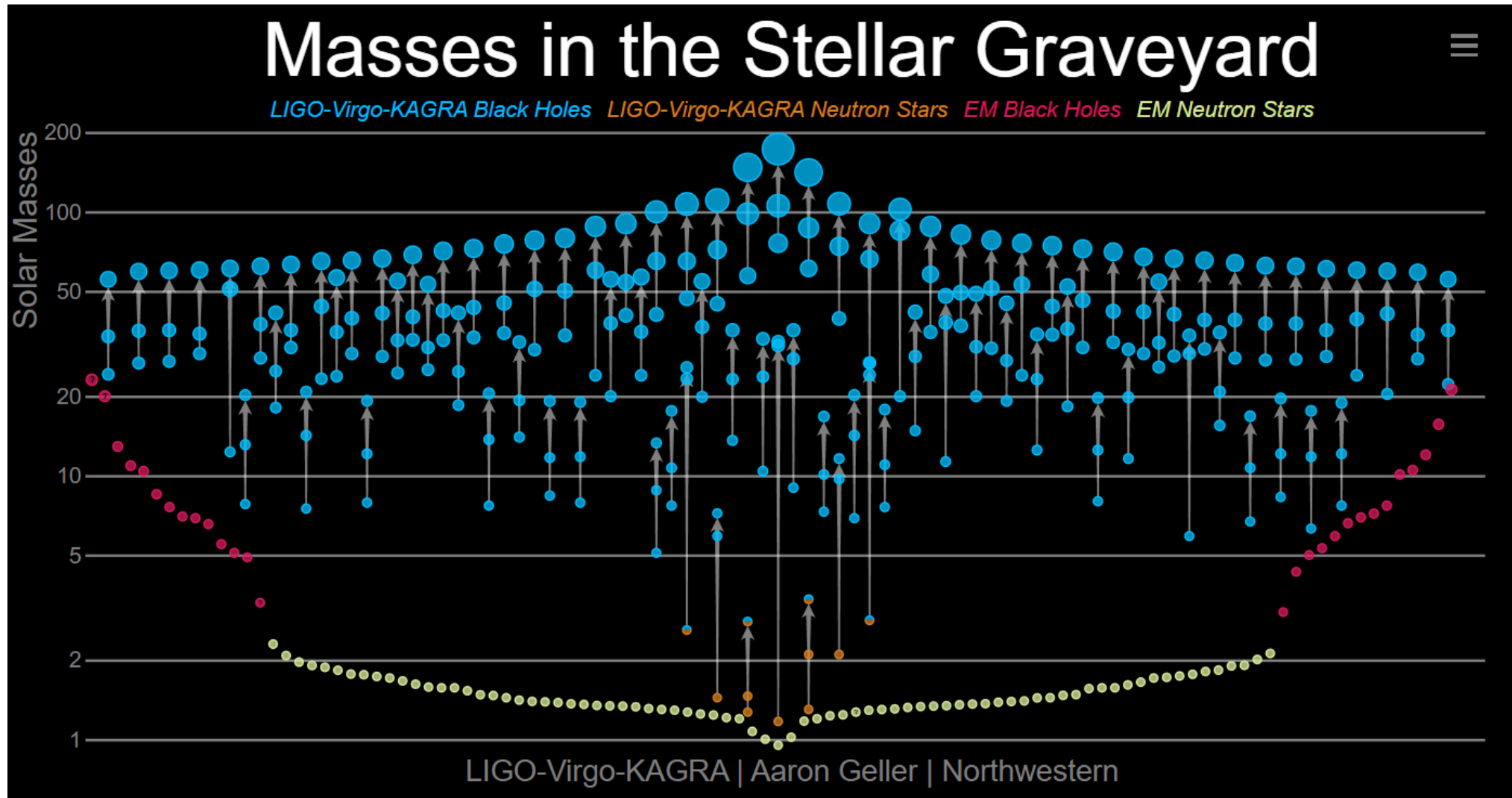


# Outline

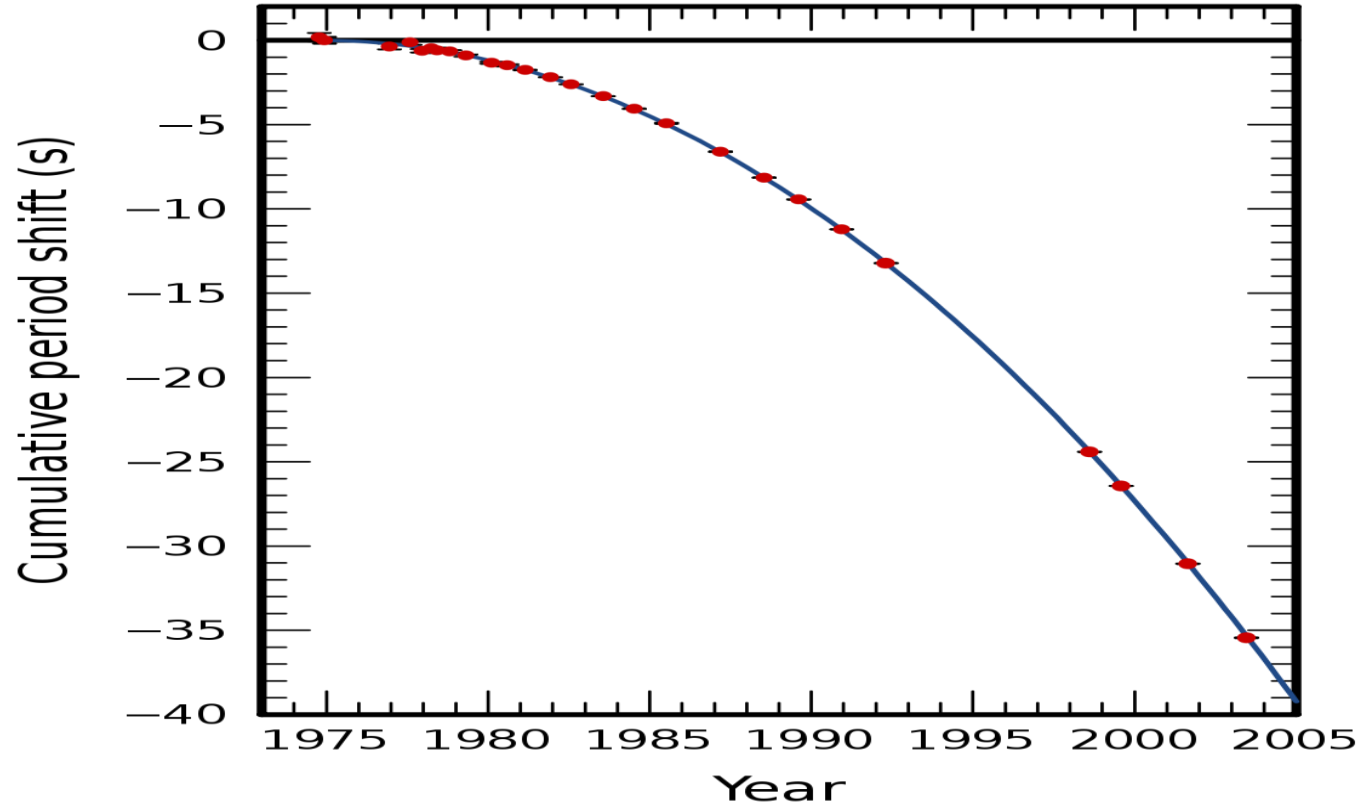
- **General motivations**
- **Physical setups**
- **Binaries in environments**
- **Final remarks**

# **General motivations**

# Main targets for GW detectors



# Time passes, and the shift tells



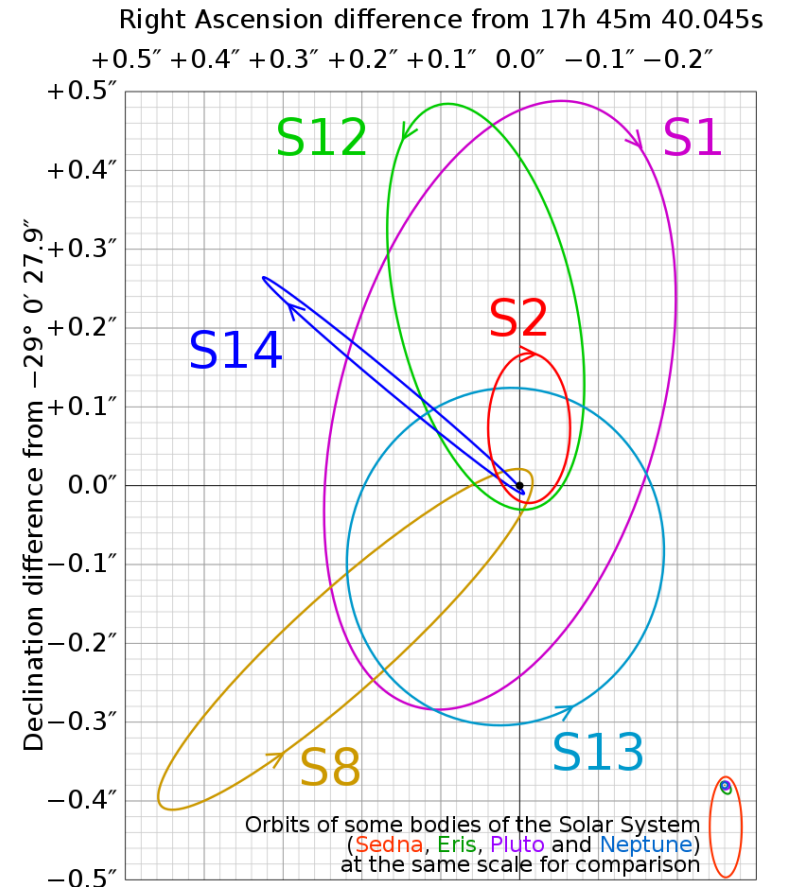
$$\frac{\dot{T}}{T} = \frac{3}{2} \frac{\text{Energy Lost}}{\text{Orbital energy}}$$

# Accretion disks, plasmas environments and fuzzy dark matter

Barausse et. al PRD **89**, 104059 (2014), Wayne Hu, Rennan Barkana, and Andrei Gruzinov Phys. Rev. Lett. 85, 1158 (2000)

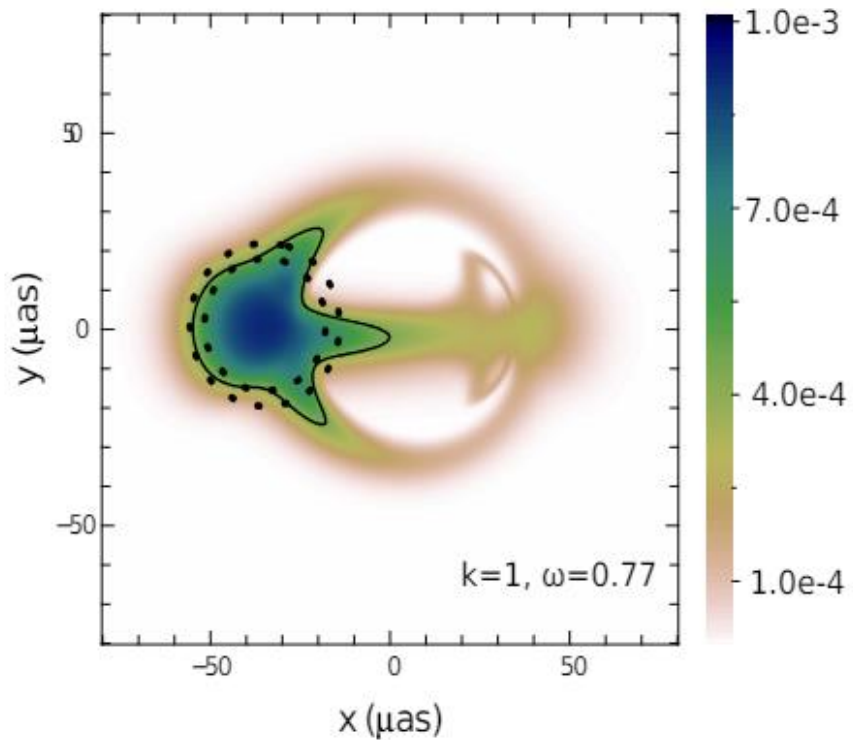


*Environmental forces  
play an IMPORTANT role*

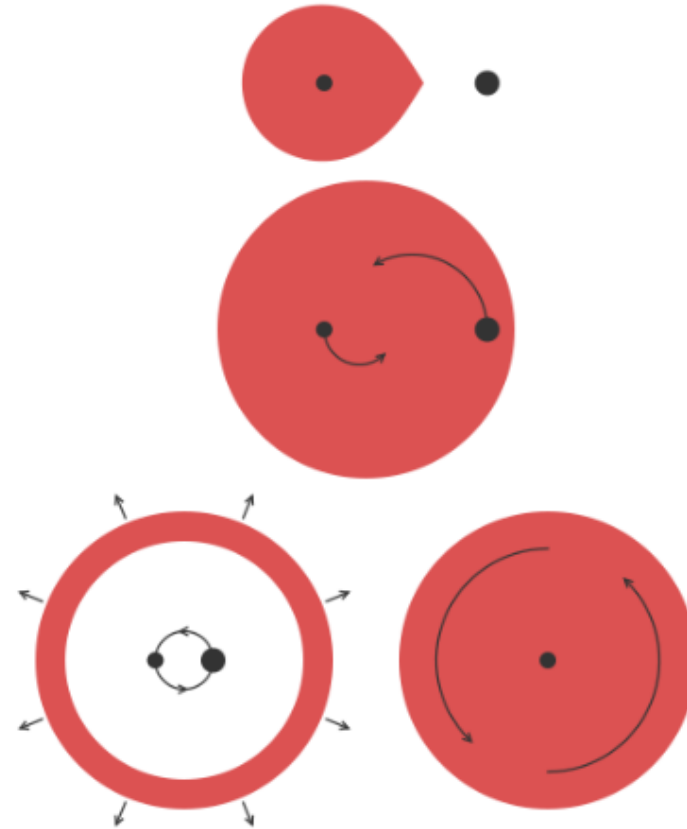


GRAVITY Collaboration

# Boson stars and common envelopes



Vicent et al. *Class.Quant.Grav.* 33 (2016) 10, 105015



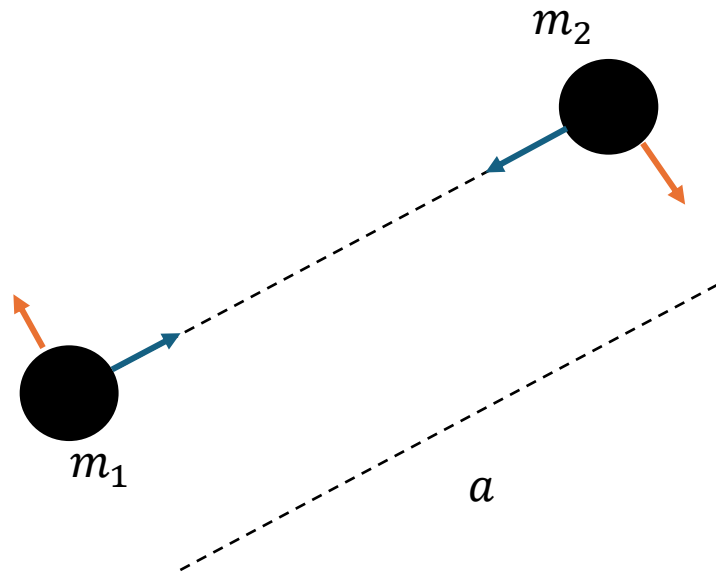
Ivanova et al., *AARev.*, 21, 59 (2013)

# Physical Setups



# Physical setups

## Newtonian elements



$$\ddot{\mathbf{r}} = f_1 \dot{\mathbf{r}} + f_2 \dot{\mathbf{R}} + f_3 \mathbf{r}$$

$$\ddot{\mathbf{R}} = f_4 \dot{\mathbf{r}} + f_5 \dot{\mathbf{R}} + f_6 \mathbf{r}$$

$$\mathbf{F}_{d,i} = -G^2 m_i^2 \rho I_d(v_i) \dot{\mathbf{r}}_i$$

$$f_1 = -\frac{G^2 M q \rho (I_{a1} + I_{a2} + I_{d1} + I_{d2})}{(q+1)^2},$$

$$f_2 = \frac{G^2 M \rho [I_{a1} + I_{d1} - q(I_{a2} + I_{d2})]}{q+1},$$

$$f_3 = GM \left\{ \frac{G^3 M q \rho^2 (I_{a1} - q I_{a2}) [I_{a1} + I_{d1} - q(I_{a2} + I_{d2})]}{(q+1)^4} - \frac{1}{r^3} \right\},$$

$$f_4 = \frac{G^2 M q \rho [q(I_{a2} - I_{d2}) - I_{a1} + I_{d1}]}{(q+1)^3},$$

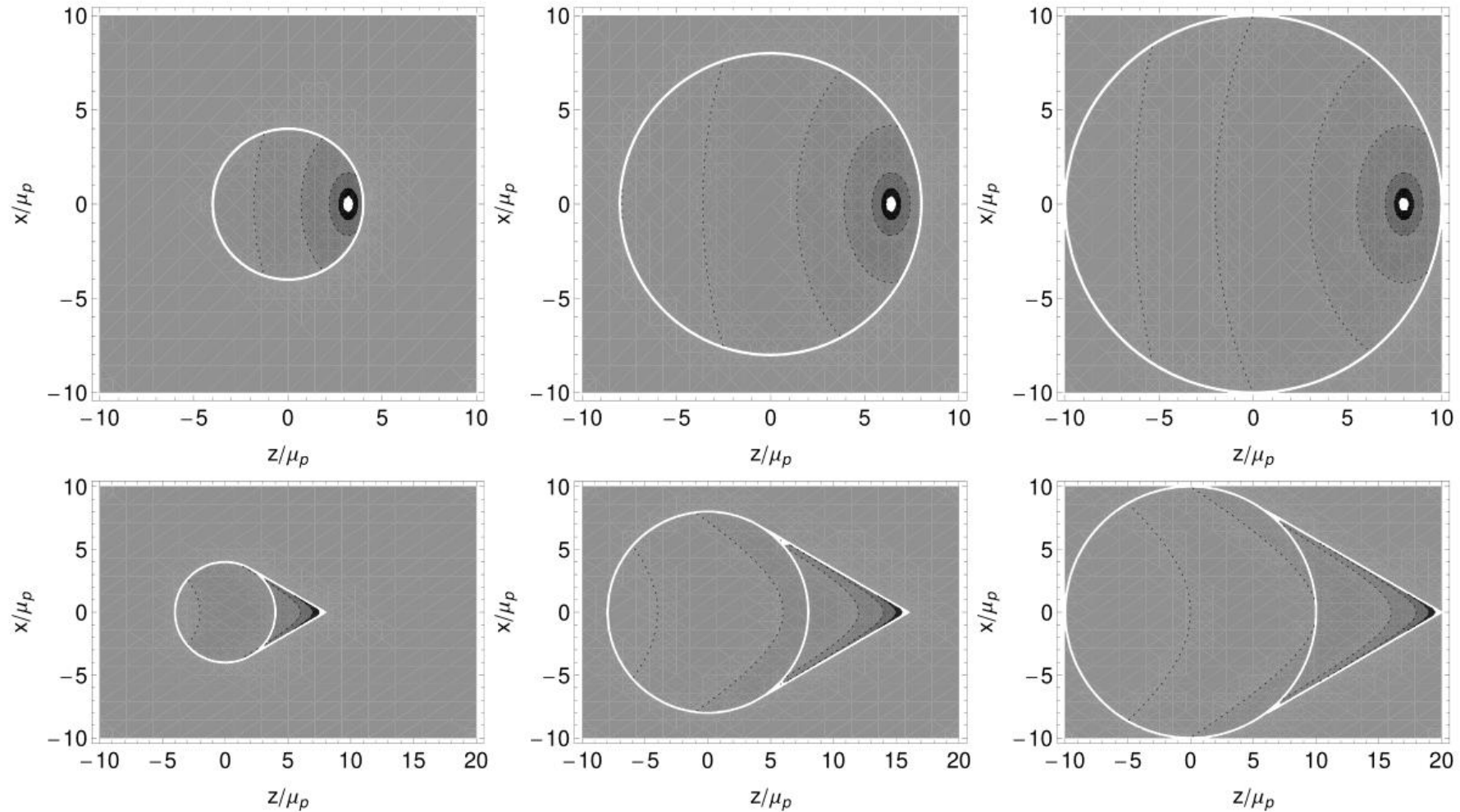
$$f_5 = -\frac{G^2 M \rho [q^2(I_{a2} + I_{d2}) + I_{a1} + I_{d1}]}{(q+1)^2},$$

$$f_6 = -\frac{G^4 M^2 q \rho^2 (I_{a1} - q I_{a2}) [q^2(I_{a2} + I_{d2}) + 2q(I_{a1} + I_{a2}) + I_{a1} + I_{d1}]}{(q+1)^5}.$$

# Physical setups

## Newtonian elements: Dynamical friction

(Chandrasekhar, APJ (1943)), (Ostriker, astro-ph/9810324 (1998))

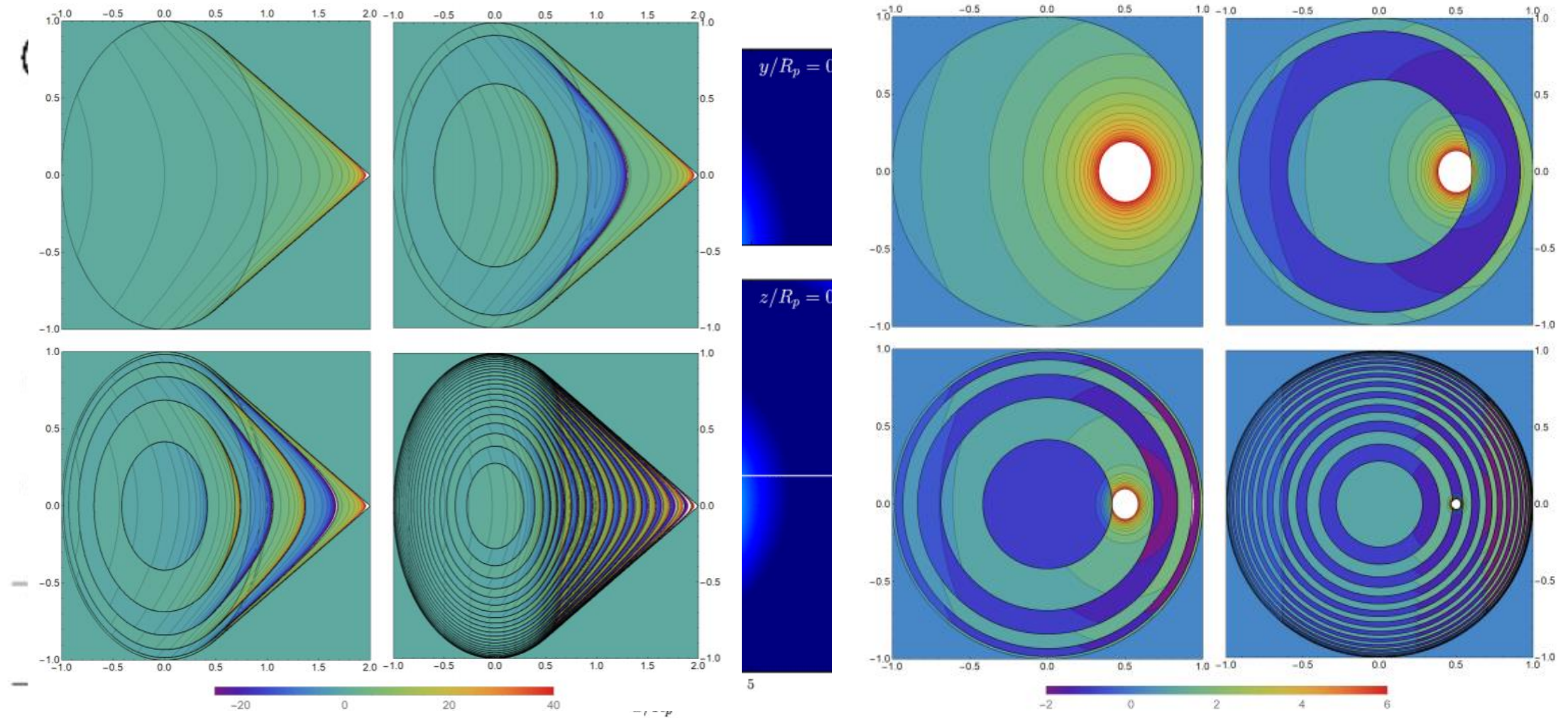


# Physical setups

## Newtonian elements: Dynamical friction in other scenarios

Kim&Kim, 0705.0084 (2007); Kim et al. 0804.2010 (2018); Vicente et. al 1905.06353 (2019)

“These are my principles. If you don’t like them, I have others”.



# Physical setups

## GR elements: Dynamical friction and scattering amplitudes

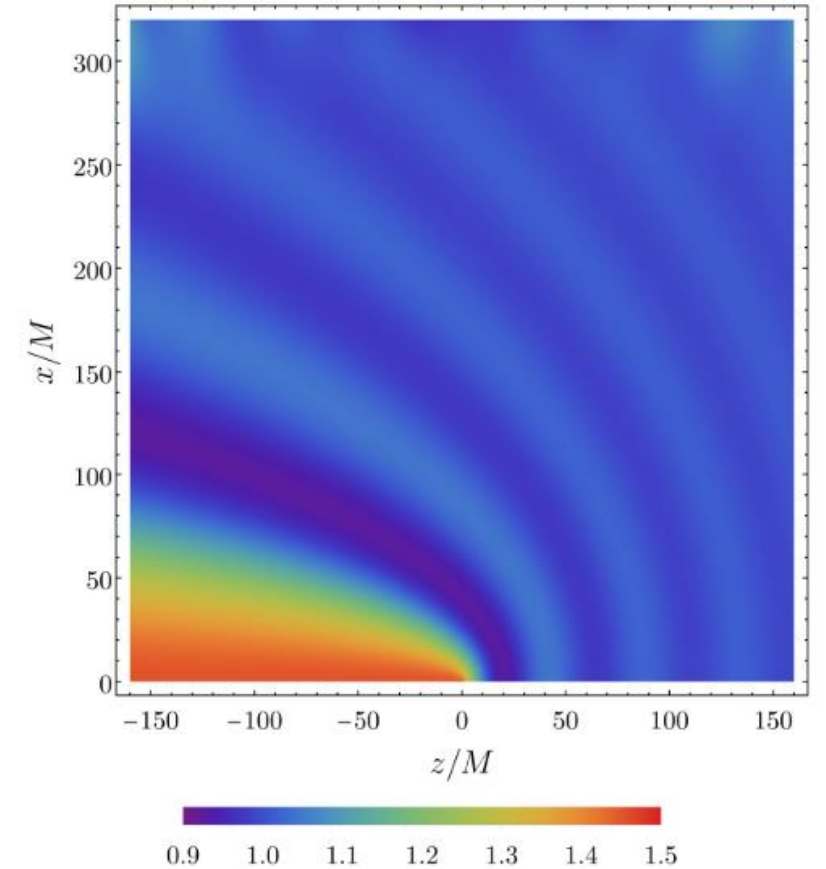
Traykova et al. *Phys. Rev. D* 104, 103014 (2021), Vicente and Cardoso. *Phys.Rev.D* 105 8, 083008 (2022), Traykova et al., *Phys.Rev.D* 108 12, L121502 (2023)

$$\dot{E}_{\text{BH}} = \frac{\pi \hbar \omega n}{\mu k_{\infty}} \sum_{\ell, m} (2\ell + 1) \frac{(\ell - m)!}{(\ell + m)!} (\text{Ps}_{\ell}^m)^2 \left( 1 - \left| \frac{R}{I} \right|^2 \right).$$

$$P_S^i(t') = \int_{S_{t'}} dV_3 T^{\alpha i} N_{\alpha}.$$



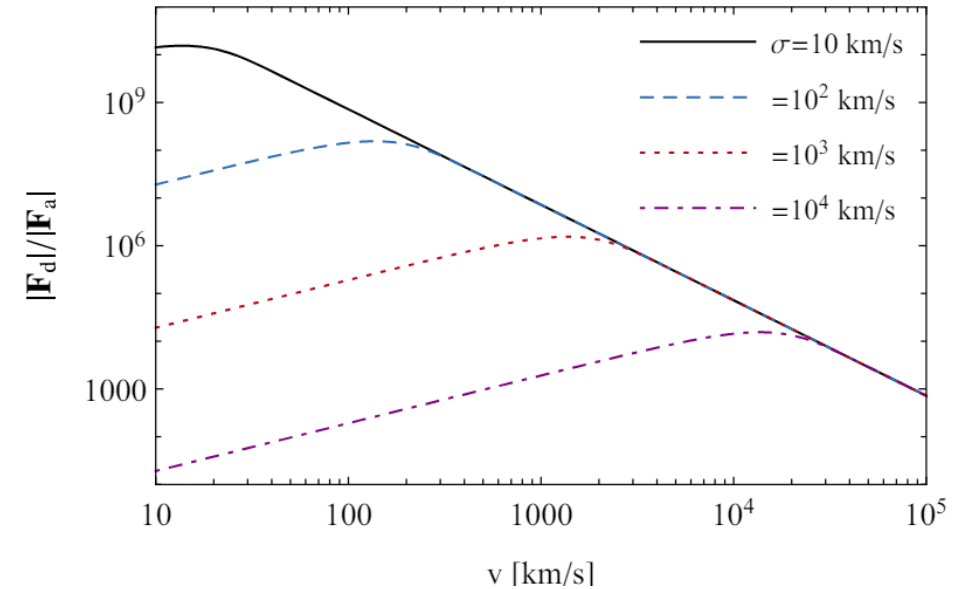
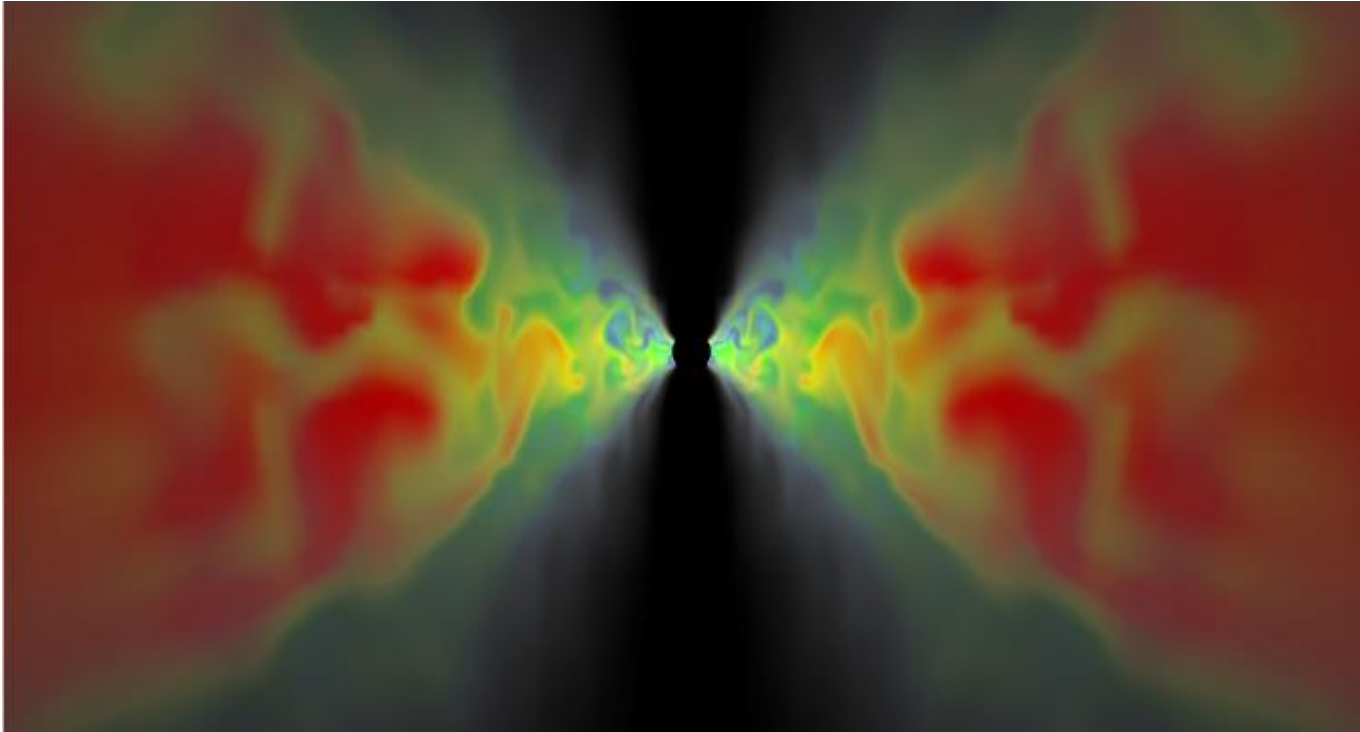
$$F' = -\frac{4\pi M^2 \rho v}{v^3} \log \left( \sqrt{1 + \frac{b_{\text{max}}^2}{(M/v^2)^2}} \right)$$



Note that the Chandrasekhar case is recovered

# Physical setups

Newtonian(-ish) elements: Accretion



$$m_i \ddot{\mathbf{r}}_i + \dot{m}_i \dot{\mathbf{r}}_i = \pm \frac{Gm_1 m_2}{r^3} \mathbf{r} + \mathbf{F}_i$$

E.g. Bondi-Hoyle-Littleton

$$\dot{m}_i = 4\pi G^2 \rho \frac{m_i^2}{(v_i^2 + c_s^2)^{3/2}}$$

# Physical setups

(Post-)Newtonian elements: quadru and octupole emissions

For **gravitational waves**, we have

$$\langle \dot{E} \rangle = -\frac{32}{5} \frac{G^4 m_1^2 m_2^2 M}{a^5 (1-e^2)^{7/2}} \left( 1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right)$$

$$\langle \dot{L} \rangle = -\frac{32}{5} \frac{G^{7/2} m_1^2 m_2^2 M^{1/2}}{a^{7/2} (1-e^2)^2} \left( 1 + \frac{7}{8} e^2 \right).$$



$$\frac{da}{dt} = -\frac{64}{5} \frac{G^3 \mu M^2}{c^5 a^3 (1-e^2)^{7/2}} \left( 1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right), \quad \downarrow$$

$$\frac{de}{dt} = -\frac{304}{15} \frac{G^3 \mu M^2}{c^5 a^4 (1-e^2)^{5/2}} \left( 1 + \frac{121}{304} e^2 \right). \quad \downarrow$$

Emission of gravitational waves  
**circularizes** orbits.

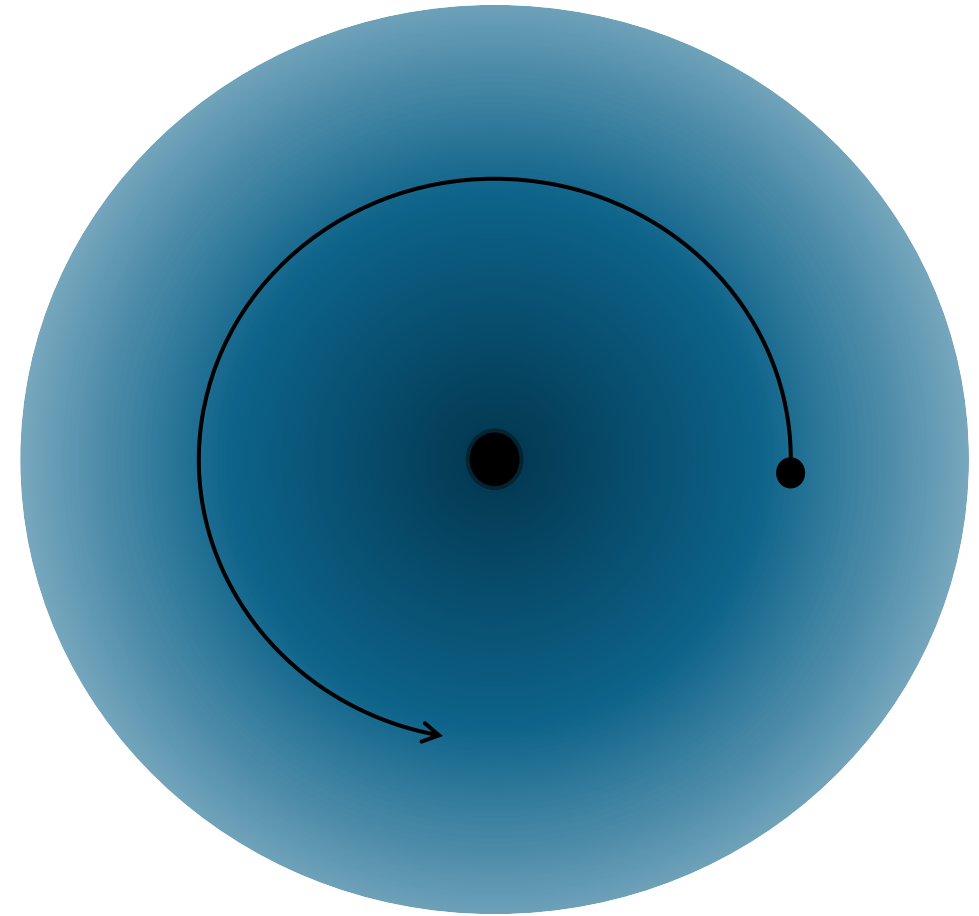
# Physical setups

Radiation: scalar fields, black holes and particles

Macedo et al. 2013, Duque et al. 2023

$$G_{\mu\nu} = 8\pi T_{\mu\nu}, \quad \square_g \Phi = \frac{\partial V}{\partial \Phi^*}$$

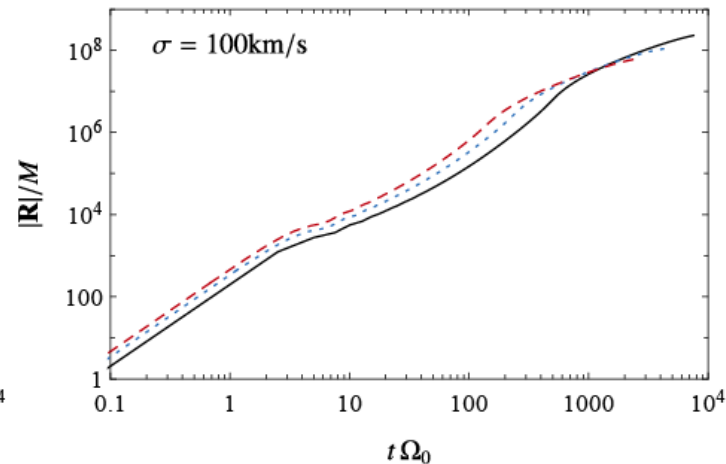
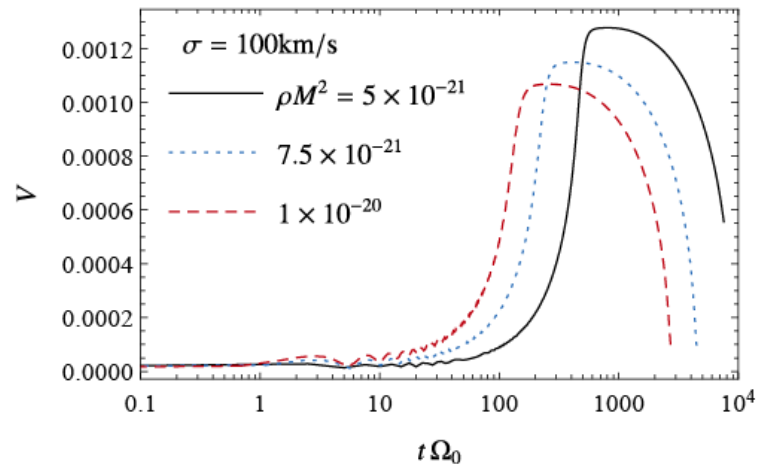
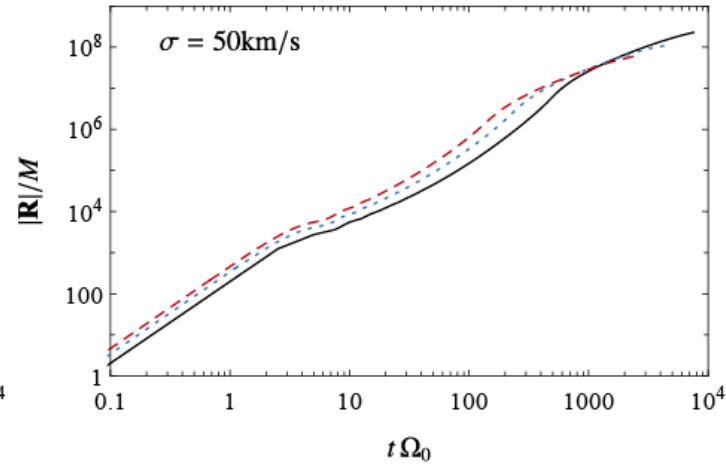
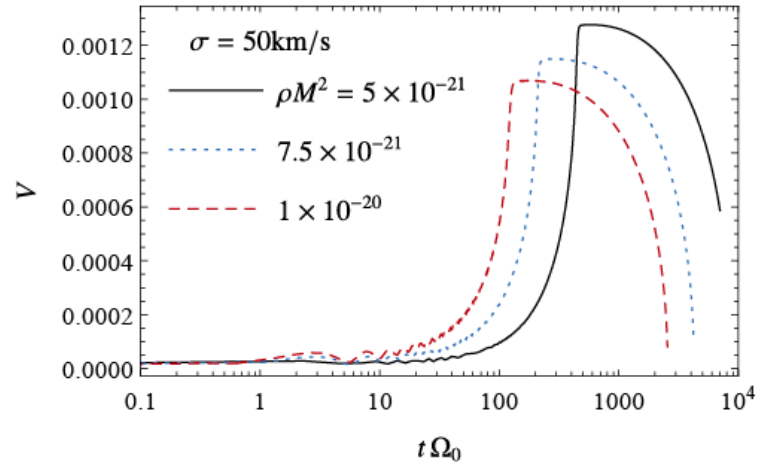
$$g_{\mu\nu} \approx \hat{g}_{\mu\nu} + q \delta g_{\mu\nu}, \quad \Phi \approx \hat{\Phi} + q \delta \Phi$$



# **Binaries in environments**



# Boost in the center of mass



$V \sim 350\text{km/s}$

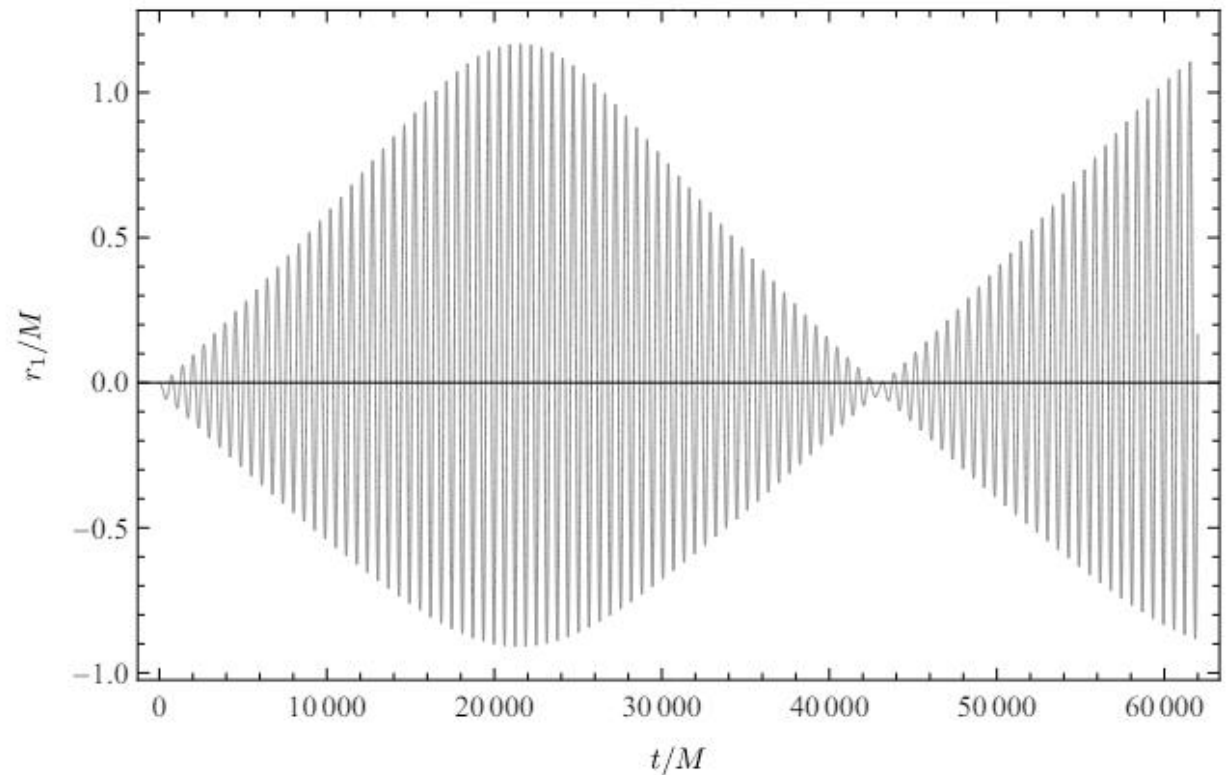
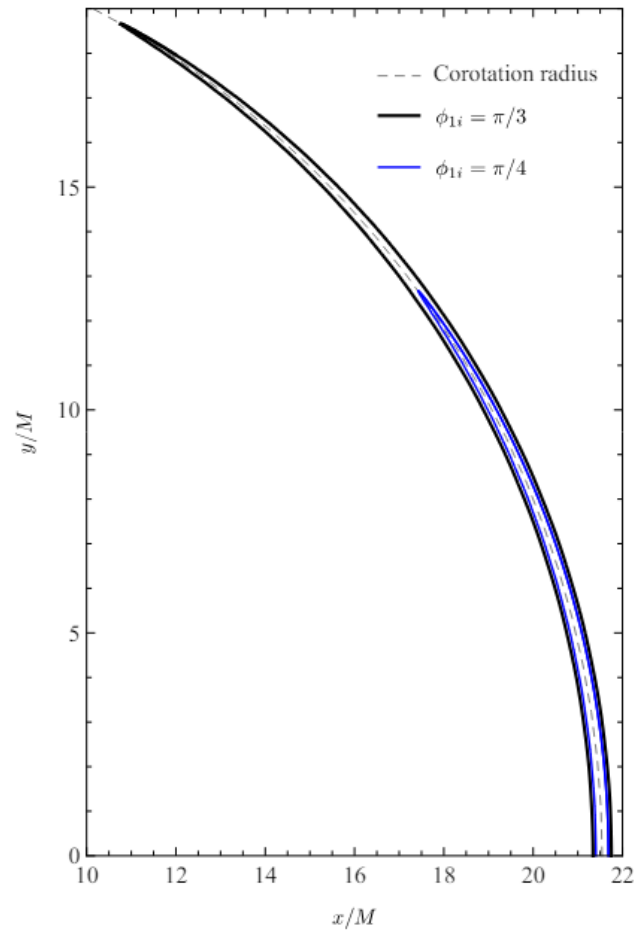
***Binaries could escape galaxies!***

$$T_{\text{max}} = 6 \left( \frac{M_{\odot}}{M} \right)^3 \left( \frac{10^{-6} \rho_{\text{water}}}{\rho} \right)^2 \text{ years}$$

# Orbital dynamics

Ferreira et al.

A time dependent background idynamics introduces richer orbital dynamics



# Eccentricity evolution and dynamical friction

Cardoso, CFBM, Vicente 2010.1515 (2021)) (Roedig&Sesana 1111.3742 (2012))

Considering dominant terms

$$\left. \begin{aligned} \langle \dot{a} \rangle &= -k\rho \sqrt{\frac{G a^5}{M}} \left( 1 + \frac{3e^2}{4} + \mathcal{O}(e^4) \right), \\ \langle \dot{e} \rangle &= \frac{3}{2} k\rho \sqrt{\frac{G a^3}{M}} e \left( 1 + \frac{3e^2}{8} + \mathcal{O}(e^4) \right). \end{aligned} \right\} \frac{da}{de} = -\frac{2a}{3e} \left( 1 + \frac{3}{8}e^2 + \mathcal{O}(e^3) \right)$$

Therefore, the environment favors eccentric motion! Considering both GW and the environment

$$\frac{da}{de} = \frac{6a \left( 5c^5 k\rho \sqrt{GM a^{11}} + 32G^3 M^4 \right)}{e \left( 304G^3 M^4 - 45c^5 k\rho \sqrt{GM a^{11}} \right)}$$



**Critical distance!**

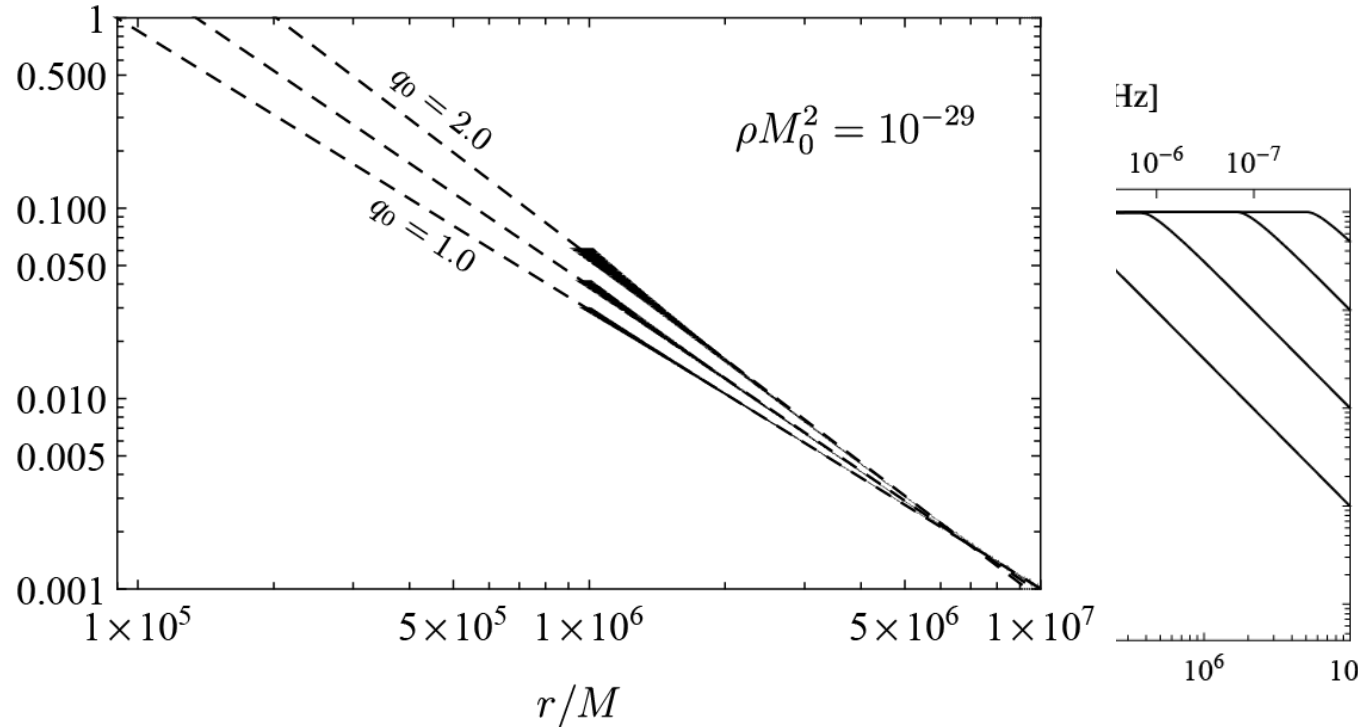
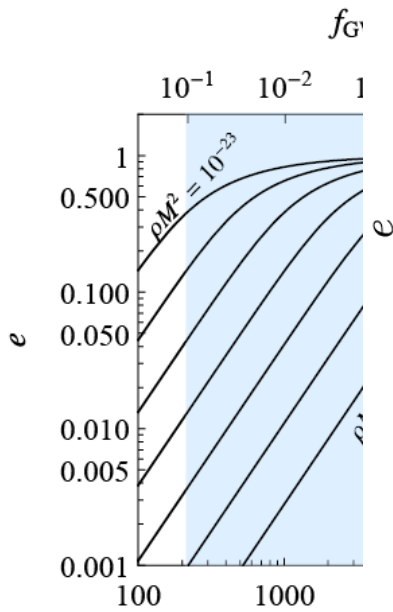
# Eccentricity evolution and dynamical friction

Cardoso, CFBM, Vicente 2010.1515 (2021)) (Roedig&Sesana 1111.3742 (2012))

$$\frac{a_c}{\left(\frac{100GM_\odot}{c^2}\right)} = 3 \times 10^4 k^{-2/11} \left(\frac{M}{100M_\odot}\right)^{7/11} \left(\frac{\rho_{10}}{\rho}\right)^{2/11}$$

$$\rho_{10} = 10^{-10} \text{g cm}^{-3}$$

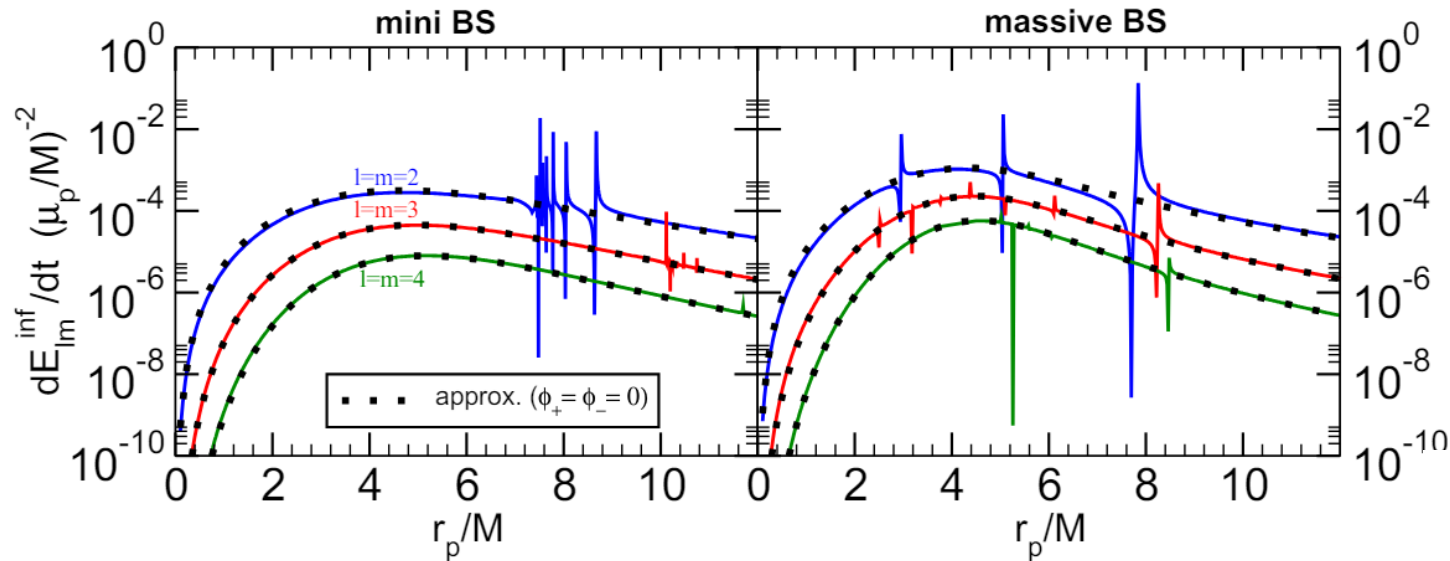
What about the evolution?



$$\frac{G^3}{c^6} \rho M^2 = 1.6 \times 10^{-24} \frac{\rho}{\rho_{10}} \left(\frac{M}{100M_\odot}\right)^2$$

# Mode excitation and dephasing

Macedo et al. Phys.Rev.D 88 6, 064046 (2013) & Astrophys.J. 774 (2013).



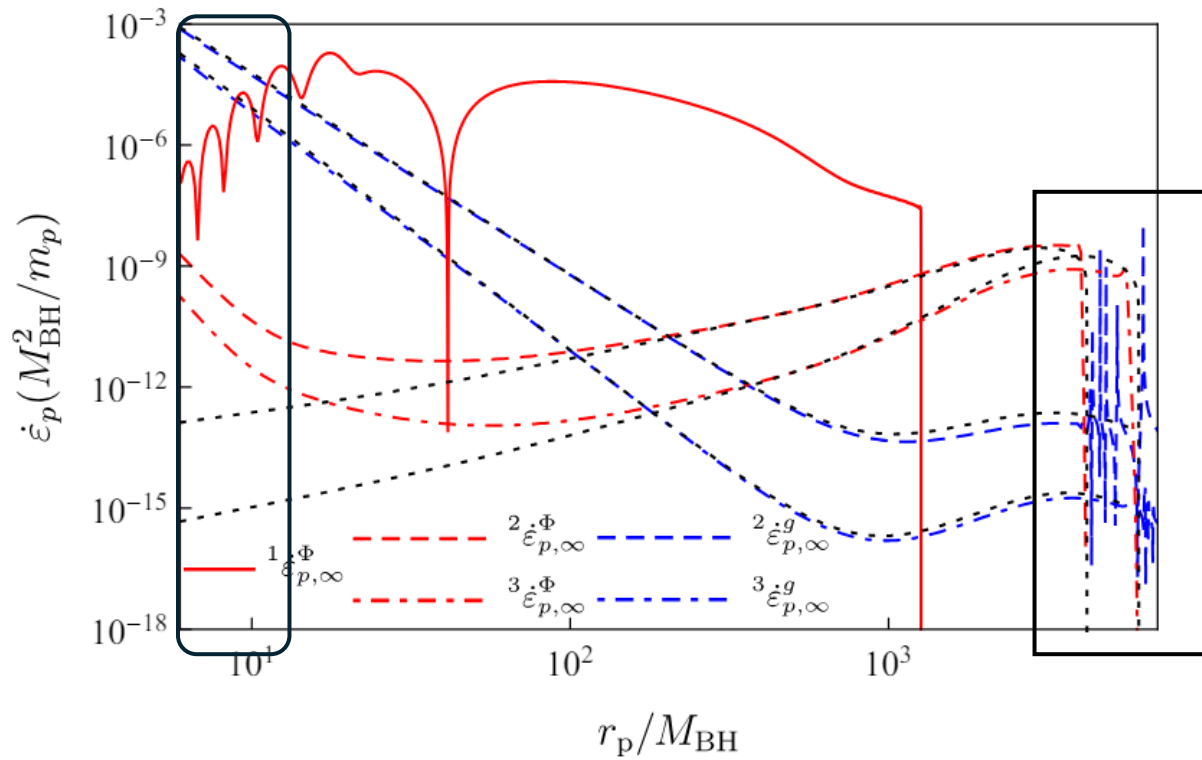
$$\dot{E} = \dot{E}_0 + h(t)\mathcal{H} [\delta t_{\text{res}}^2 - (t - t_{\text{res}})^2]$$

$$|\delta\phi_{\text{GW}}| \sim \frac{5h}{16\nu^2} \frac{M\Delta\Omega_{\text{res}}}{(M\Omega_{\text{res}})^{10/3}} \frac{T_{\text{obs}}}{M}$$

$$|\delta\phi_{\text{GW}}| \approx 8.6 \times 10^3 \text{rads} \left[ \frac{10^5 M_{\odot}}{M} \right] \left[ \frac{T_{\text{obs}}}{1\text{yr}} \right] \\ \times \left[ \frac{h/\nu^2}{10^{-2}} \right] \left[ \frac{\sigma_I M}{10^{-5}} \right] \left[ \frac{0.1}{\sigma_R M} \right]^{10/3}$$

# EMRI: BH surrounded by (scalar) dark matter

Duque et al. arXiv:2312.06767 (2023)



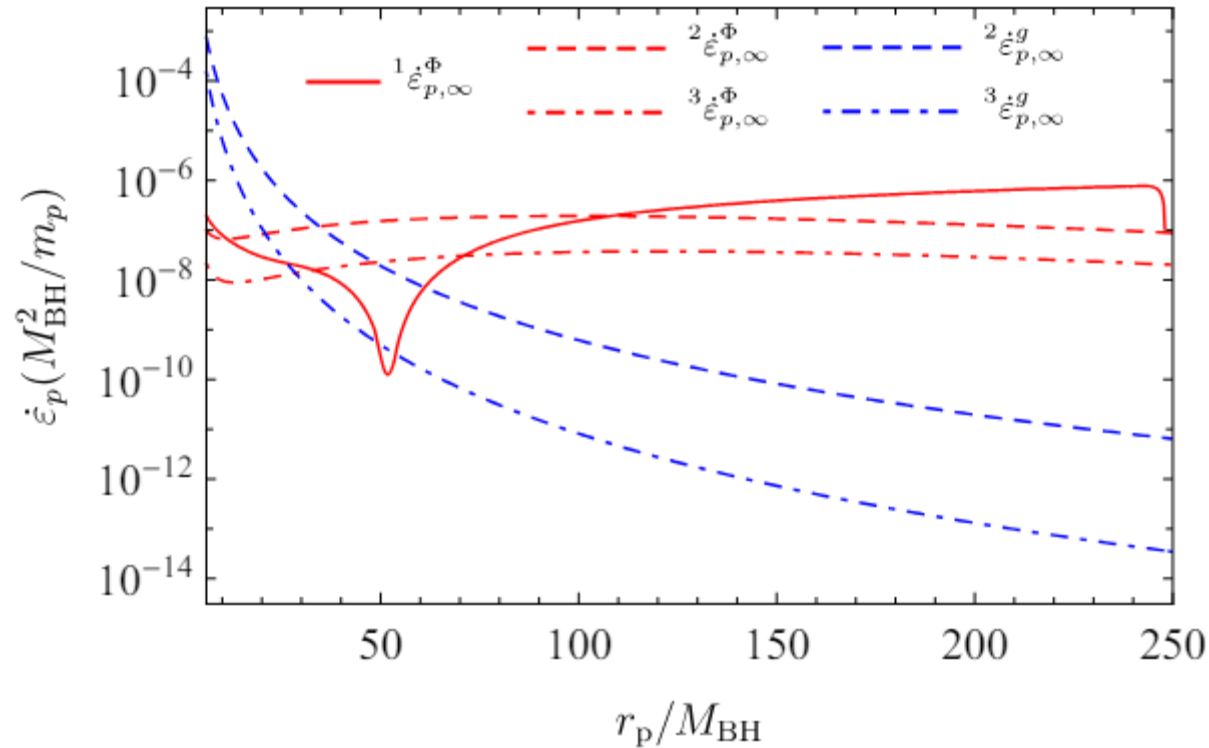
$$\frac{\dot{T}}{T} = \frac{3}{2} \frac{\text{Energy Lost}}{\text{Orbital energy}}$$

For Sgr A\* like scenarios

$$\delta_{\text{cyc}} \sim \left( \frac{t_{\text{obs}}}{t_{\Phi}} \right) \#_{\text{cyc}} \gtrsim 200 \text{ cycles}$$

# EMRI: BH surrounded by a cloud

Duque et al. arXiv:2312.06767 (2023). Brito, Shah Phys.Rev.D 108 8, 084019 (2023)



$$\delta_{\text{cyc}} \sim 20 \left( \frac{q}{10^{-5}} \right) \left( \frac{t_{\text{obs}}}{4 \text{ yr}} \right) \left( \frac{\Omega_p/\pi}{0.3 \text{ mHz}} \right)^{1/3}$$

# Final remarks

- Binaries in environments can be remarkably different.
- Dynamical friction is important in eccentricity and CM boosts.
- Period shifts can be a discriminator.
- Axisymmetric time-dependent environment introduces resonances.



Thank you!



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